LANDAU COLLISIONS IN THE PARTICLE BASIS WITH PETSCPIC

PETSc 2023

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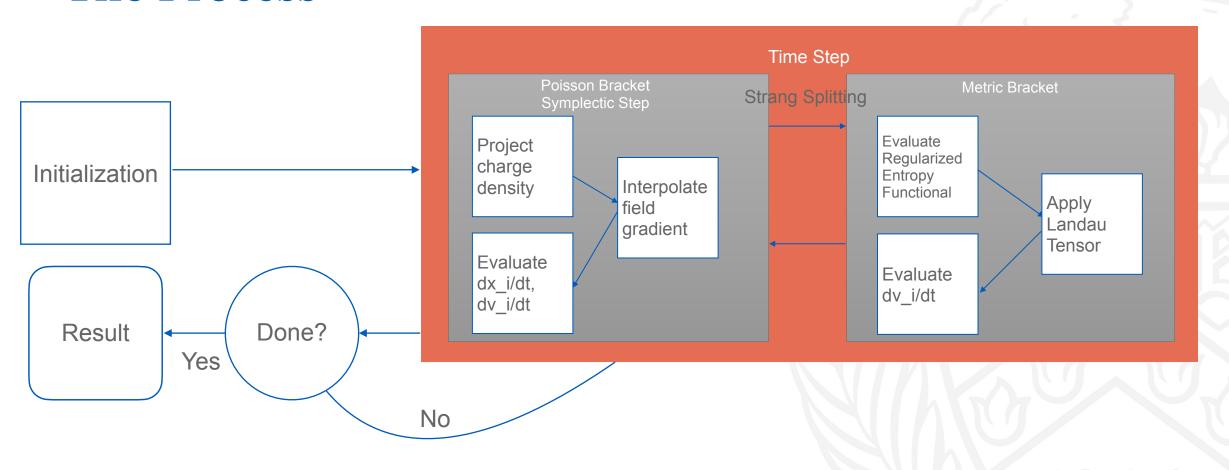




Metriplectic Formulation

$$\frac{dU}{dt} = (U, F) + \{U, F\}$$

The Process



Landau Collision Integral

- Represent transfer of momentum in the grazing limit
 - Non relativistic particles
- Conservative in the continuum
 - Conservative FE discretization in PETSc by Adams et al
 - PETSc implmentation in plexlandau/ex1 and ex2
 - Particle discretizations remained challenging

$$\begin{split} \frac{df_{\alpha}}{dt} &= \sum_{\beta} C_{\alpha\beta} [f_{\alpha}, f_{\beta}] \\ C_{\alpha\beta} &= \nu_{\alpha\beta} \frac{m_0}{m\alpha} \nabla \cdot \int_{\bar{\Omega}} d\bar{v} U(v, \bar{v}) \cdot (\frac{m_0}{m_\alpha} \bar{f}_{\beta} \nabla f_{\alpha} - \frac{m_0}{m_\beta} f_{\alpha} \bar{\nabla} \bar{f}_{\beta}) \\ U(\vec{v}, \bar{v}) &= \frac{1}{|\vec{v} - \bar{v}|^3} (|\vec{v} - \bar{v}|^2 I - (\vec{v} - \bar{v})(\vec{v} - \bar{v})) \end{split}$$

Grid Approach

- FE implementation by Adams et al
- High order elements, non conforming adaptive mesh refinement
- Conservative
- Works at exascale
- Well verified (stationary Maxwellian, effective reproduction of Spitzer model of plasma friction)
- Supports multiple species per grid

Adams, Wang, Merson, Knepley (2022). A performance portable, fully implicit Landau collision operator with batched linear solvers

Adams, Hirvijoki, Knepley, Brown, Isaac, Mills (2017). Landau Collision Integral Solver with Adaptive Mesh Refinement on Emerging Architectures

Adams, Brennan, Knepley, Wang (2022) Landau collision operator in the CUDA programming model applied to thermal quench plasmas

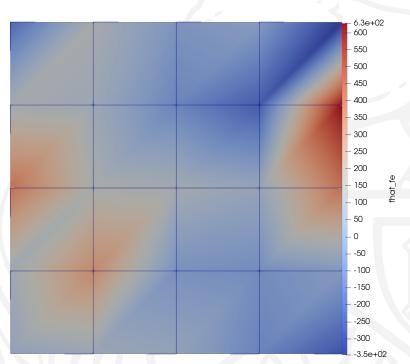


Fig: Arbitrary field projected to the grid from the particle basis

$$f_{FE} = \sum_{i} \phi_{i} \hat{f}$$

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Basis Projection

- Incompatible representations
- Compare continuous functions with radon measure
- Need to transition between representations
- Enforce weak equivalence

$$\int_{\Omega} \phi_i \sum_j f_j \phi_j = \int_{\Omega} \phi_i \sum_p w_p \delta(\hat{x} - \hat{x}_p)$$

$$\sum_j f_j \int_{\Omega} \phi_i \phi_j = \sum_p w_p \int_{\Omega} \phi_i \delta(\hat{x} - \hat{x}_p) \longrightarrow M\hat{f} = M_P \hat{w}$$

$$\int f_{FE} = \int f_p$$

$$\int v f_{FE} = \int v f_p$$

$$\int_{\Omega} \phi_i f_{FE} = \int_{\Omega} \phi_i f_P \quad \forall \phi_i \in \mathcal{T}$$

Projection Operator

- Particle-Mesh interaction required
- Deposit particles to grid to calculate charge density/potential fields
- Synthesize particles from mesh to evaluate changes in weighting
- Deposition Mass Matrix inversion (easy!)
- Synthesis Pseudo Inverse of particle mass matrix (not as easy:()
 - KSPLSQR (no PC, heavy)
 - preconditioned KSPLSQR w/ normal equations and ASM/icc or BJacobi/LU

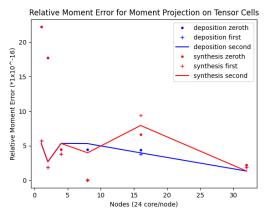
$$\int_{\Omega} \phi_i \sum_j \hat{f}_j \phi_j = \int_{\Omega} \phi_i \sum_p \hat{w}_p \delta(x - x_p)$$

$$\sum_{j} \hat{f}_{j} \int_{\Omega} \phi_{i} \phi_{j} = \sum_{p} \hat{w}_{p} \int_{\Omega} \phi_{i} \delta(x - x_{p})$$

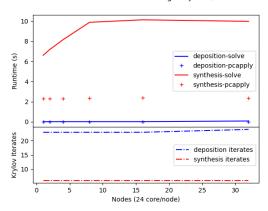
$$M\hat{f} = V\hat{w}$$

$$\hat{f} = M^{-1} M_P \hat{w} \qquad \qquad \hat{w} = M_P^{\dagger} M \hat{f}$$

Projection Operator

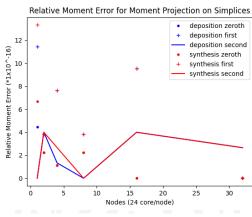


Tensor Element Weak Scaling of Bjacobi/LU

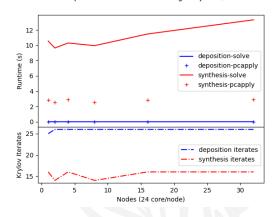


Relative Error





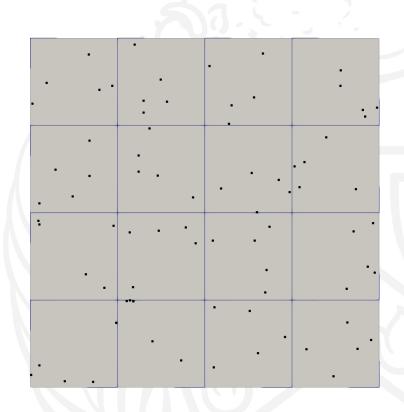
Simplicial Element Weak Scaling of Bjacobi/LU



Particle Approach

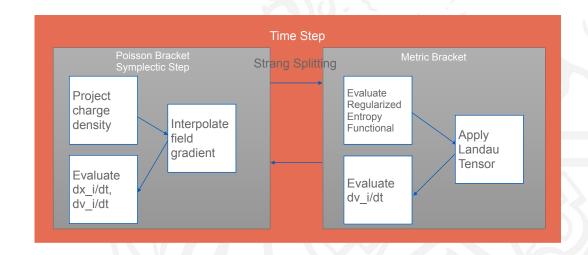
• Say we have no velocity grid...

$$f_P = \sum_i w_i \delta(x - x_p)$$



Metriplectic Formulation

- Notion of Metriplectic structure, two brackets
 - Poisson, Metric
- Demonstrably conserve symplectic structure in the Poisson bracket
- Metric bracket guarantees Hamiltonian as a Casimir of the bracket (mass, momentum, energy conservation)
- Monotonic entropy production



$$\frac{dU}{dt} = (U, F) + \{U, F\}$$

$$C_{\alpha\beta} = \nu_{\alpha\beta} \frac{m_0}{m\alpha} \nabla \cdot \int_{\bar{\Omega}} d\bar{v} U(v, \bar{v}) \cdot \left(\frac{m_0}{m_\alpha} \bar{f}_\beta \nabla f_\alpha - \frac{m_0}{m_\beta} f_\alpha \bar{\nabla} \bar{f}_\beta\right)$$

Collisions in the Particle Basis

- Carillo et. al demonstrate Maxwellian steady state of the particle representation of the collision operator
 - Interpret Landau collision integral as gradient flow over set of probability measures
 - Cast in terms of entropy functional
 - velocity evolution determined as variational derivative of the entropy functional
 - Regularize Entropy
 - Mollification procedure

$$f_p = \sum_p w_p \delta(v - v_p)$$

$$S = \int f lnf$$

$$U_{\epsilon} = -\int Q(\xi) (\nabla_{\nu} \frac{\delta S}{\delta f} - \nabla_{\bar{\nu}} \frac{\delta S}{\delta \bar{f}}) d\bar{\nu}$$

^[1] Carillo et al, A particle method for the homogeneous Landau equation

^[2] Hirvijoki, Structure-preserving marker-particle discretizations of Coulomb collisions for particle-in-cell codes

^[3] Zonta et al, Multispecies structure-preserving particle discretization of the Landau collision operator

Mollification

- Integrate distribution against smooth function
 - Gaussians
- Essentially perform a convolution with smooth set of functions
 - Generate smooth, differentiable representation of the basis of delta functions
- This is not carried out in application
 - Just replace delta function with gaussian at particle coordinate

$$\frac{dv_p}{dt} = \sum_{\bar{p}} \nu_{p\bar{p}} 1(p,\bar{p}) Q(\xi) \Gamma(S_{\epsilon}, p, \bar{p})$$

$$\psi_{\epsilon} = \frac{1}{2\pi\epsilon} e^{\frac{-(v^2)^2}{2\epsilon}}$$

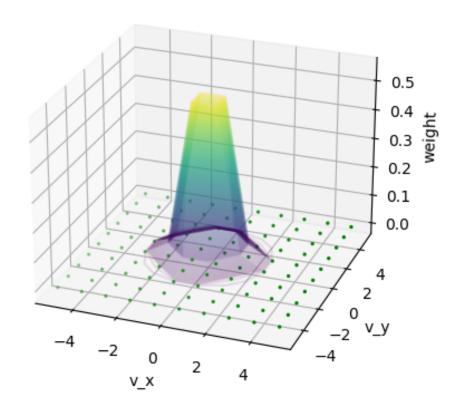
$$S \to S_{\epsilon} = -\int \psi * f \ln \psi * f$$

$$S = -\int_{\Omega} \sum_{p} w_{p} \psi_{\epsilon}(v - v_{p}) ln \sum_{k} w_{k} \psi_{\epsilon}(v - v_{k})$$

$$\Gamma(S_{\epsilon}, p, \bar{p}) = \frac{\partial S}{m \partial v_p} - \frac{\partial S}{m \partial v_{\bar{p}}}$$

Initialization

- Convergence of the method only guaranteed for sufficiently regular particle spacing
 - · Lay particles out in a grid
 - Calculate weight based on initial conditions in that region of velocity space
- Regularization parameter limits particle irregularity
 - Prohibits QMC methods (ie Sobol sequence Quiet Starts)
 - Uniformly weighted particles under investigation
 - regularization parameter in terms of discrepency of the set



Stepping

- Precompute entropy functional gradients for each species
- Form the Jacobian
 - Approximated either Finite Difference or Matrix Free
 - SNESFD (Heavy)
 - SNESMF (Less heavy)
- Compute residual over all pairs
- Copy back to Swarm

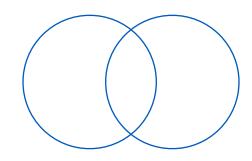
$$\frac{v^{n+1} - v^n}{\Delta t} = \sum_{\bar{p}} 1(p, \bar{p}) Q(\xi^{n+1/2}) \Gamma(S^n, p, \bar{p})$$

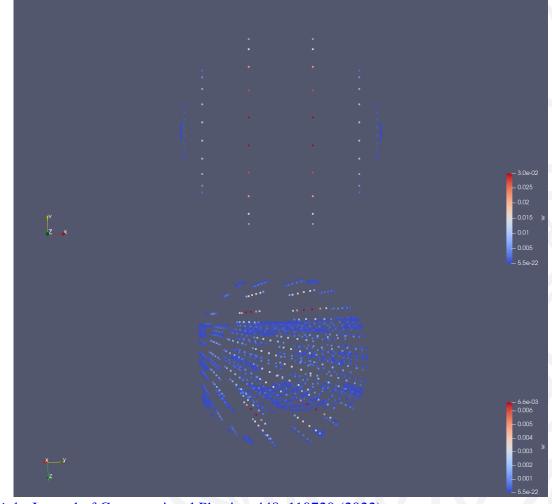
Entropy Functional Gradient

$$\frac{\partial S_{\epsilon}}{\partial v_p} = -\int \nabla_{v_p} \sum_{p} w_p \psi_{\epsilon}(v - v_p) \ln \sum_{k} w_k \psi_{\epsilon}(v - v_k)$$
$$= -\int w_p \psi'(v - v_p) (1 + \ln \sum_{k} w_k \psi(v - v_k))$$

Evaluation of the Entropy Functional Gradient

- Gauss-Hermite Quadrature
- AlgoIM High order quadrature over implicitly defined domains
 - Can generate high order quadrature of complex domains including cusps and intersections of arbitrary polynomial functions
 - Trimmed quadrature space can be leveraged to increase speedup

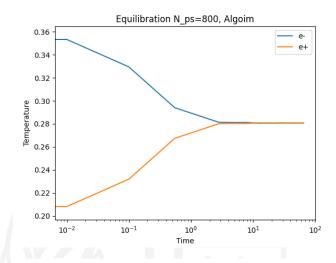


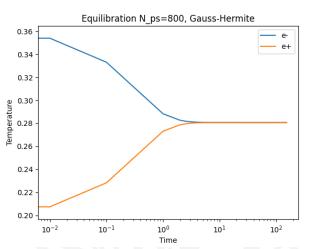


Particle Basis Landau

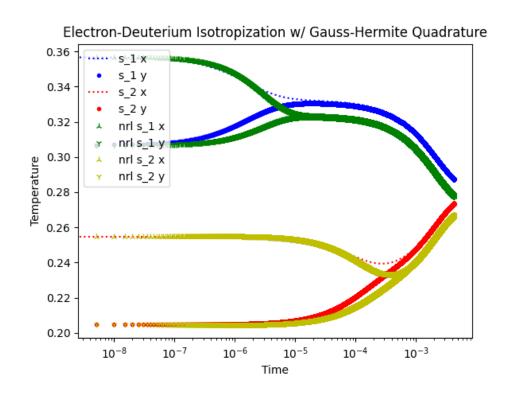
- Recent theoretical work makes evaluation of the integral in the particle basis possible with proof of conservative properties
- Initial implementation non due to time discretization conservative
- Hirvijoki proposed new time discretization
- Implicit midpoint like time stepping
 - TSTHETA w/ theta = 0.5

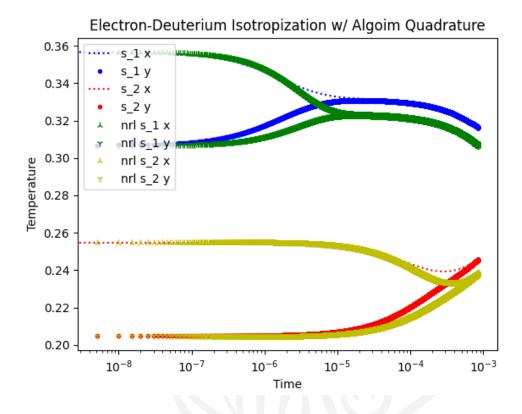
$$\frac{v^{n+1} - v^n}{\Delta t} = \sum_{\bar{p}} 1(p, \bar{p}) Q(\xi^{n+1/2}) \Gamma(S^n, p, \bar{p})$$





Isotropization





Conclusion

- Conservative implementation from Adams et all in FE basis
 - Highly optimized
 - Multispecies support
- Conservative implementation in particle basis implementation underway
 - Undergoing optimization process
 - Multispecies support
- Monolithic system is a subject of future development
 - Relaxation RK, Monolithic Discrete Gradients



Regularized Entropy Form

- Cast as advective
- Particle basis entropy functional non smooth/continuous
- Mollify with smooth RBF

$$\begin{split} \frac{\partial f}{\partial t} &= \frac{\partial}{\partial v} \cdot \int_{\Omega} Q(v - \bar{v}) \{ f(\bar{v}) \frac{\partial f}{\partial v} - f(v) \frac{\partial f}{\partial \bar{v}} \} \\ \partial_{v} f &= f \partial_{v} ln f = -f \partial_{v} \frac{\delta S}{\delta f} \\ S &= -\int_{\Omega} f ln f \\ \frac{\partial f}{\partial t} + \frac{\partial}{\partial v} \cdot (Ff) &= 0 \\ F(v_{p}) &= \sum_{\bar{v}} w_{\bar{p}} Q(v_{p} - v_{\bar{p}}) \{ \frac{\partial}{\partial v_{p}} \frac{\delta S}{\delta f} - \frac{\partial}{\partial v_{\bar{p}}} \frac{\partial S}{\partial f} \} \end{split}$$

$$\psi_{\epsilon} = \frac{1}{2\pi\epsilon} e^{\frac{-(v^2)}{2\epsilon}}$$

$$f * \psi = \int_{\Omega} \psi(v) f(v) dv$$

$$S[f * \psi_{\epsilon}] = -\int_{\Omega} f * \psi_{\epsilon} lnf * \psi_{\epsilon}$$